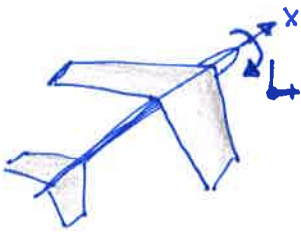


Lesson 19

Lateral S+C  
"roll"

# Roll Moment



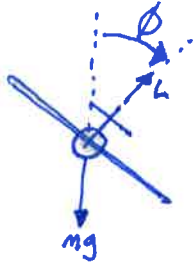
$$C_L = \frac{L}{\rho S b}$$

Warning

L is a moment in this context

L is a lift force in other contexts.

Is there static roll stability?

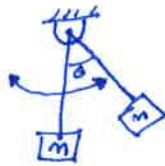


In other words, does  $\phi$  decrease to zero?

No. Roll angle is statically neutral.

Q: Is there a pendulum effect?

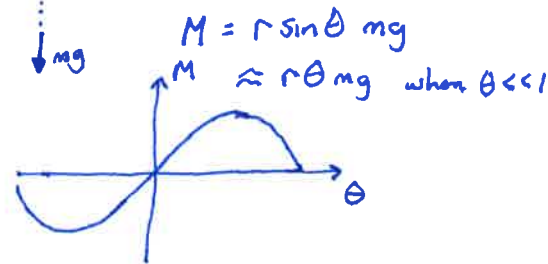
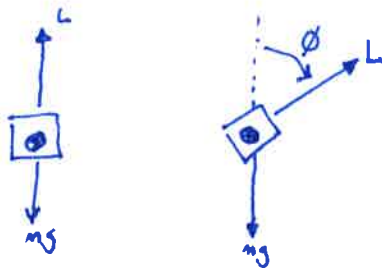
No.



About the rotation point (of a pendulum),

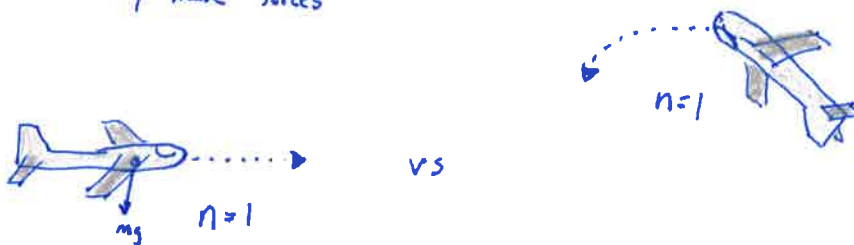


Q: Why is the aircraft different?

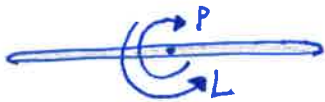


In ~~steady~~ flight, the roll angle does not impact the lift. ~~the lift is independent of the roll angle~~

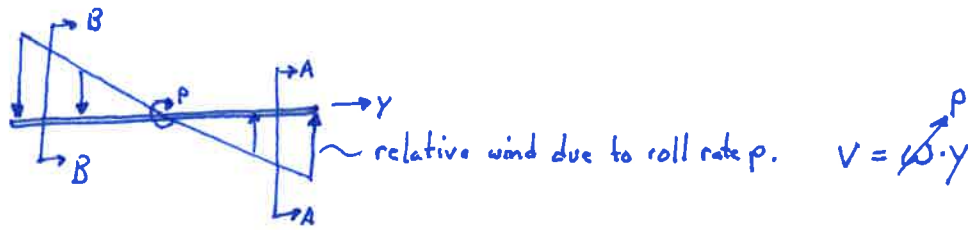
The aircraft (and the pilot) generate aerodynamic force regardless of the direction of Earth. In fact, in flight you can't know down or up by the local body frame forces



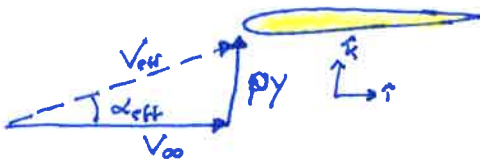
# Roll Rate Damping (Nelson 3.6.4) This is a dynamic term.....



A right rolling wing generates a left (-) roll moment.



At section A-A:



$$\bar{V}_{eff} = \bar{V}_{\infty} \hat{i} + p y \hat{k}$$

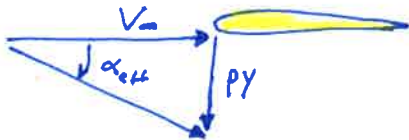
$$\alpha_{eff} = \alpha \tan\left(\frac{p y}{V_{\infty}}\right) \approx \frac{p y}{V_{\infty}} \text{ when } p y \ll V_{\infty}$$

This section will generate a positive lift due to the roll rate

$$\frac{\Delta L}{\Delta y} \approx -g \bar{c} C_{L_{\alpha}} \alpha \approx -g \bar{c} C_{L_{\alpha}} \frac{p y}{V_{\infty}}$$

Assume

At section B-B

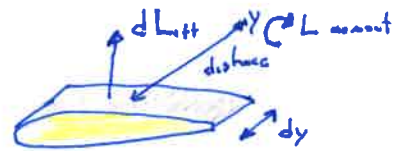


$$\alpha_{eff} \approx \frac{p y}{V_{\infty}} \text{ since } y < 0 \text{ then } \alpha_{eff} < 0$$

This section generates negative lift due to roll rate

Roll moment due to roll rate (Strip theory)

$$\frac{dL}{dy} = -\text{Lift} \cdot \text{distance} = -\left(g \bar{c} C_{L_{\alpha}} \frac{p y}{V_{\infty}}\right) \cdot (y)$$



Integrate over the entire span

$$\int dL = \int_{-b/2}^{+b/2} -g \bar{c} C_{L_{\alpha}} \frac{p y}{V_{\infty}} \cdot y dy \Rightarrow L = \int_{-b/2}^{+b/2} -g \bar{c} C_{L_{\alpha}} \frac{p y^2}{V_{\infty}} dy$$

Non-dimensionalize  $C_L = \frac{L}{g S b}$

$$C_L = -\frac{1}{S b} \int_{-b/2}^{+b/2} \bar{c}(y) C_{L_{\alpha}} \frac{p y^2}{V_{\infty}} dy$$

Constant Chord wing  $\bar{c}(y) = \bar{c}$

$$C_L = -\frac{\bar{c} C_{L\alpha} \rho}{5b V_\infty} \int_{-b/2}^{b/2} y^2 dy = -\frac{\bar{c} C_{L\alpha} \rho}{5b V_\infty} \left. \frac{y^3}{3} \right|_{-b/2}^{b/2} = C \cdot \frac{b^3}{8} \left( -\frac{b^3}{8} \right) \frac{b^3}{4}$$

$$= -\frac{\bar{c} C_{L\alpha} \rho}{5b V_\infty} \frac{b^3}{4}$$

Dynamic Simulation:

$$IC \begin{cases} \phi(t=0) = 0 \\ \dot{\phi}(t=0) = 180^\circ/s \end{cases}$$



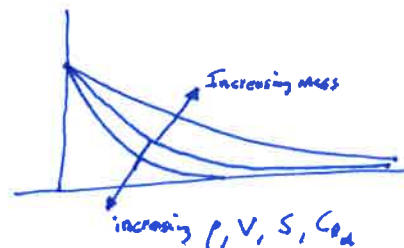
$$\begin{cases} \text{Gov} \\ \text{Eqs} \\ \text{ID} \end{cases} \begin{cases} I \ddot{\phi} = L = \rho S b C_L \\ \ddot{\phi} = \frac{\rho S b}{I} C_L \end{cases}$$

$$= \frac{\rho S b}{I} (-1) \frac{\bar{c} C_{L\alpha} \rho}{5b V_\infty} \frac{b^3}{4}$$

$$\ddot{\phi} = -\frac{\rho}{V_\infty} \cdot \bar{c} \cdot \frac{4}{M b^2} \frac{b^3}{4} C_{L\alpha} \dot{\phi} = -\frac{\rho}{V_\infty} \bar{c} \frac{b}{M} C_{L\alpha} \dot{\phi}$$

$$\text{Thus, } \dot{r} = -\frac{1}{2} \rho V \frac{\bar{c} b}{M} C_{L\alpha} r = -\frac{1}{2} \rho V \frac{S}{M} C_{L\alpha} r$$

$$\text{Solution to } \dot{r} = -C r \text{ is } r = r_0 e^{-Ct} = 180^\circ/s e^{-\frac{1}{2} \rho V \frac{S}{M} C_{L\alpha} t}$$



The wing damping drives the rate to zero.

$$r \rightarrow 0 \text{ as } t \rightarrow \infty$$

Mass:  $M$

Span:  $b$

$I$ : moment of inertia

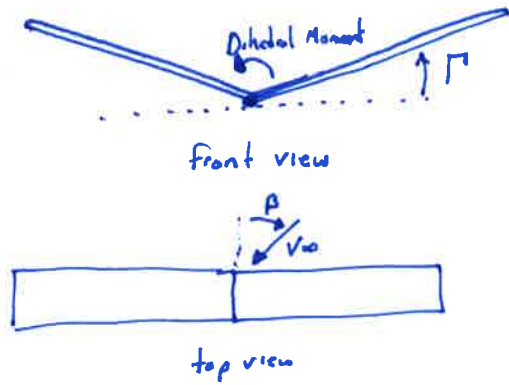
$$= \int r^2 dm$$

$$= \int_{-b/2}^{b/2} r^2 \rho A dr$$

$$= \rho A \frac{b^3}{4} = \frac{M}{b} \frac{b^3}{4}$$

$$= \frac{M b^2}{4}$$

# Dihedral

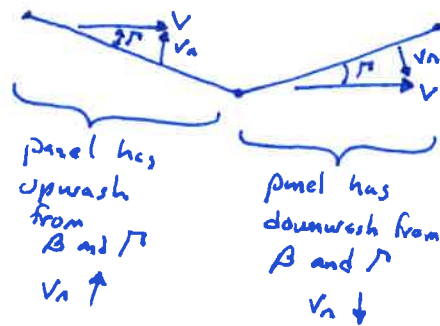
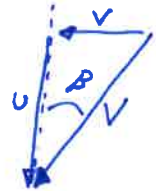


Dihedral <sup>creates</sup> gives a roll moment given a side slip angle  $\beta$ .

## Simplified Analysis

$$\beta \approx a \sin\left(\frac{v}{|V|}\right) \approx \frac{v}{U} \quad \text{when } v \ll U$$

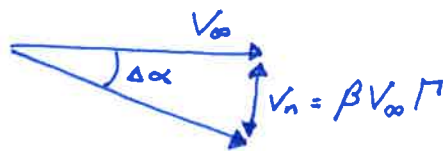
$$\text{Thus } v \approx \beta V_{\infty}$$



In the wing panel frame, there is a normal component of the side velocity  $v$ .

$$V_n = v \sin \Gamma \approx v \Gamma \\ \approx \beta V_{\infty} \Gamma$$

Now that the normal component is known, compare to  $V_{\infty}$

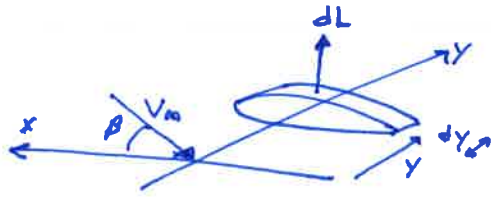


$$\tan(\Delta\alpha) = \frac{v_n}{V_{\infty}} = \frac{\beta V_{\infty} \Gamma}{V_{\infty}}$$

For small angles,  $\tan(\Delta\alpha) \approx \Delta\alpha = \beta \Gamma$

The local change in AOA is the product of sideslip angle and the dihedral angle

Integrate over the entire wing.



$$dL_{\text{moment}} = -dL_{\text{lift}} \cdot y$$

$$\uparrow$$

$$g c C_{L\alpha} \Delta \alpha dy$$

$$\begin{cases} +\beta \Gamma & \text{when } y > 0 \\ -\beta \Gamma & \text{when } y < 0 \end{cases}$$

$$dL = -2g c C_{L\alpha} \beta \Gamma y dy$$

Integrate

$$\int dL = \int_{-b/2}^{b/2} -2g c C_{L\alpha} (\pm \beta \Gamma) y dy$$

$$L = \int_0^{b/2} -2g c C_{L\alpha} \beta \Gamma y dy$$

$$C_L = \frac{1}{g S b} \int_0^{b/2} -2g c C_{L\alpha} \beta \Gamma y dy = -\frac{2C_{L\alpha} \beta}{S b} \int_0^{b/2} c \Gamma y dy$$

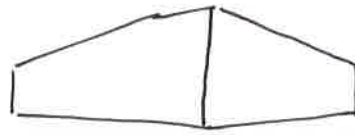
$$C_L = -\frac{2C_{L\alpha} \beta}{S b} \int_0^{b/2} c \Gamma y dy$$

$$C_{L\beta} = -\frac{2C_{L\alpha}}{S b} \int_0^{b/2} c \Gamma y dy$$

We call  $C_{L\beta}$  the dihedral effect

A positive  $\beta$  creates a negative roll moment

# Linearly Tapered Wing



$$c = c_r \left( 1 - (1-\lambda) \left| \frac{2y}{b} \right| \right)$$

$$C_{L\beta} = \frac{-2C_{L\alpha}}{Sb} \int_0^{b/2} c_r \left( 1 - (1-\lambda) \left( \frac{2y}{b} \right) \right) \Gamma y dy$$

$$= \frac{-2C_{L\alpha}}{Sb} \cdot \frac{C_r b^2 (2\lambda + 1) \Gamma}{24}$$

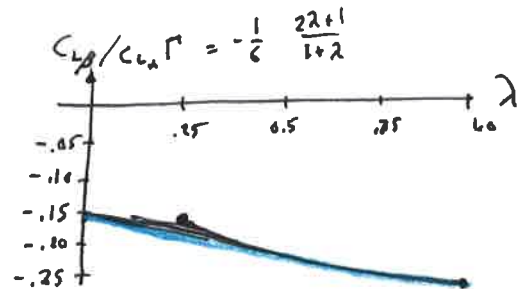
but  $AR = \frac{2b}{C_r(1+\lambda)}$  verify on own  $\Rightarrow C_r = \frac{2b}{AR(1+\lambda)}$

$$C_{L\beta} = \frac{-2C_{L\alpha}}{8b} \cdot \frac{2b}{AR(1+\lambda)} \cdot \frac{b^2}{24} (2\lambda + 1) \Gamma \cdot \left| \frac{AR^2}{b^2} \right|$$

cross off terms

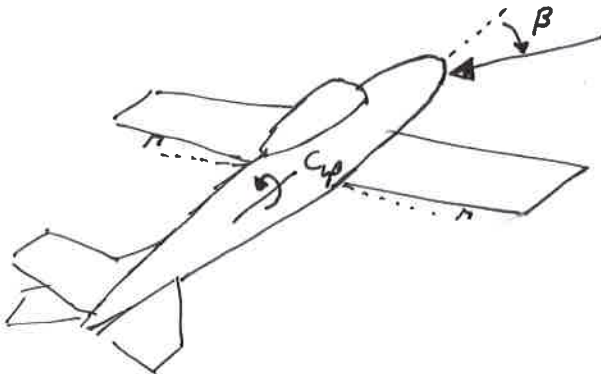
$$C_{L\beta} = -\frac{1}{6} \left( \frac{2\lambda + 1}{1 + \lambda} \right) C_{L\alpha} \Gamma$$

linear taper



The dihedral effect is linear with respect to dihedral angle

Wings with no taper ( $\lambda=1$ ) have stronger dihedral effect than  $\lambda \approx 0$ .



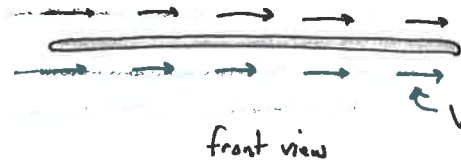
Dihedral allows the rudders to provide roll control

$$\delta_r \Rightarrow \beta = \frac{-C_{n\delta_r}}{C_{n\beta}} \delta_r \Rightarrow C_L = C_{L\beta} \beta$$

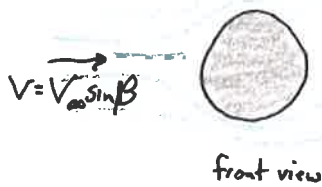
$$C_L = -\frac{C_{n\delta_r}}{C_{n\beta}} C_{L\beta} \delta_r$$

# Dihedral Effect from Fuselage

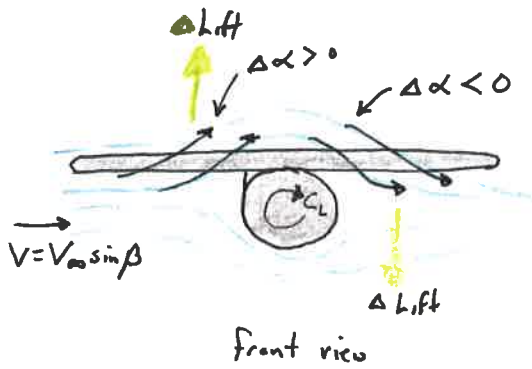
The wing position contributes to dihedral effect



No fuselage, no  $C_{L\beta}$



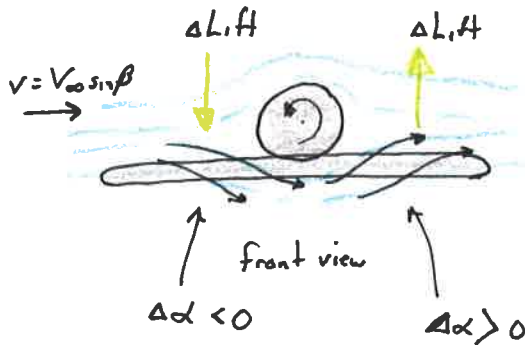
fuselage, no  $C_{L\beta}$   
perhaps  $C_{Y\beta}$  (sideforce)



fuselage + high wing

In this case,  $\beta$  from left (front view) gives a roll moment in negative x direction

$$C_{L\beta}^{\text{high wing + fuselage}} < 0$$



fuselage + low wing

$$C_{L\beta}^{\text{low wing}} > 0$$

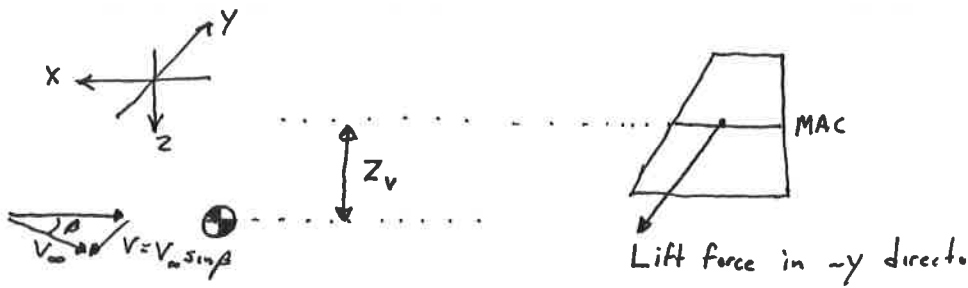
McCormack says:

$$\Delta C_{L\beta}^{\text{high}} \approx -0.00016 \frac{1}{\text{deg}}$$

$$\Delta C_{L\beta}^{\text{low}} \approx 0.00016 \frac{1}{\text{deg}}$$



# Dihedral Effect From Vertical Tail



$$L = -L_v z_v = -z_v \eta_v \rho C_{L_{\alpha v}} (\beta + \sigma) S_v$$

Non-dim

$$C_L = \frac{L}{\rho S b} = -\eta_v C_{L_{\alpha v}} \left( \frac{z_v S_v}{S b} \right) (\beta + \sigma)$$

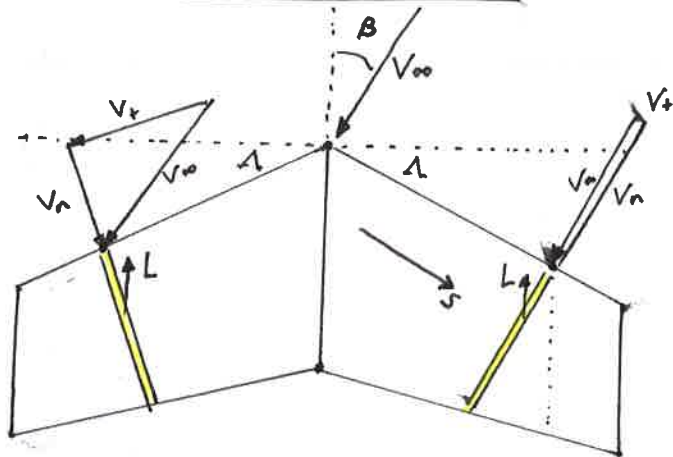
$$\text{with } \sigma = \sigma_0 + \frac{d\sigma}{d\beta}$$

$$C_{L_{\beta}} = -\eta_v C_{L_{\alpha v}} \left( \frac{z_v S_v}{S b} \right) \left( 1 + \frac{d\sigma}{d\beta} \right)$$

A vertical above the cg (common!)  
tends to provide dihedral effect.

$$C_{L_{\beta}} < 0$$

# Dihedral Effect Due to Swept Wings



The normal velocity on the left wing is  
 $V_n = V_\infty \cos(\Delta + \beta)$

The normal velocity on the right wing is  
 $V_n = V_\infty \cos(\Delta - \beta)$

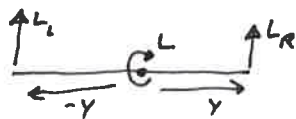
The lift is

$$L = \rho_n c_n C_L = \rho_\infty \cos^2(\Delta + \beta) \cdot \cos \Delta c C_L$$

The right wing lift is

$$L = \rho_n c_n C_L = \rho_\infty \cos^2(\Delta - \beta) \cdot \cos \Delta c C_L$$

The differential roll moment is



$$dL = \rho c C_L y (\cos^2(\Delta + \beta) - \cos^2(\Delta - \beta)) \cos \Delta \frac{dy}{\cos \Delta}$$

and ~~from~~  $dy = \cos \Delta ds$

$$dL = \rho c C_L y (\cos^2(\Delta + \beta) - \cos^2(\Delta - \beta)) dy$$

Integrate over half span (since we already took differential)

$$L = \rho c C_L (\cos^2(\Delta + \beta) - \cos^2(\Delta - \beta)) \int_0^{1/2} c y dy$$

The total lift is

$$C_{L_{\text{net}}} = C_L \cos^2 \Delta \quad (\text{approximately for swept wings})$$

The roll moment is thus

$$C_L = \frac{L}{\rho S b} = \frac{C_L}{S b \cos^2 \Delta} (\cos^2(\Delta + \beta) - \cos^2(\Delta - \beta)) \int_0^{b/2} c y dy$$

$$C_{L\beta} = \left. \frac{dC_L}{d\beta} \right|_{\beta=0} = -\frac{4C_L \tan \Delta}{S b} \int_0^{b/2} c y dy$$

Or for a linearly tapered wing

$$C_{L\beta} = -\frac{1+2\lambda}{3(1+\lambda)} C_L \tan \Delta$$

This is only an approximate result.

• The dihedral effect in swept wings depends on the lift

$$C_L \uparrow \Rightarrow C_{L\beta} \downarrow \quad \text{more dihedral effect!}$$

• " depends on  $\tan$  of sweep angle

• " stronger for larger  $\lambda$