

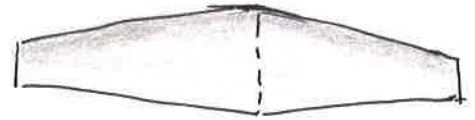
Lesson 16

Prandtl Lifting Line

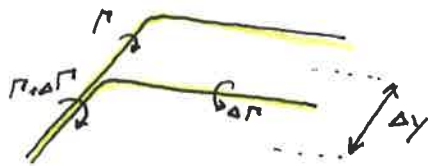
# Prandtl Lifting Line

## Assumptions:

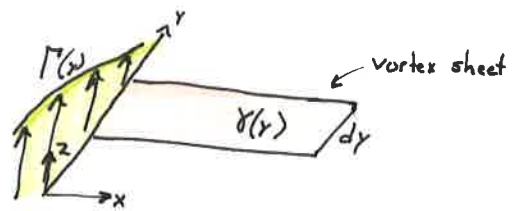
- Locally 2D chordwise flow
- No sweep
- No dihedral (flat)
- $AR > 4$
- $\alpha$  in linear  $C_{e\alpha}$  region (we will discuss another approach using experimental data!)



## Wake strength



⇒



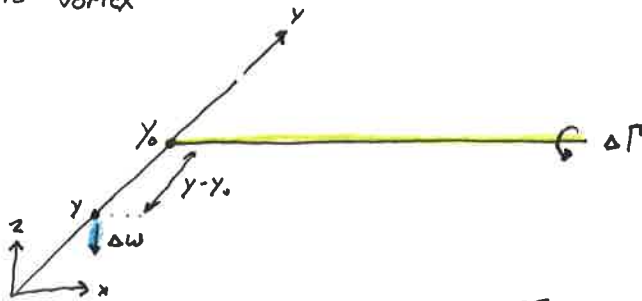
$$\frac{-(\Gamma + \Delta\Gamma) + (\Gamma)}{\Delta y} = \Delta\gamma \Rightarrow \frac{\Delta\Gamma}{\Delta y} = \Delta\gamma$$

A decrease in the bound vortex must be accompanied by an increase in the wake vortex (to satisfy Helmholtz)

$$\frac{d\Gamma}{dy} = -\gamma$$

$$\gamma = -\frac{d\Gamma}{dy}$$

## Single semi infinite vortex



What is the velocity  $\Delta w$  in the (-) z direction consistent/induced from a semi-infinite vortex of strength  $\Delta\Gamma$ ?

$$\Delta w = -\frac{\Delta\Gamma}{4\pi(y_0 - y)}$$

$$V = \frac{\Gamma}{4\pi h} \leftarrow \Delta\Gamma$$

$\Delta w$        $y - y_0$

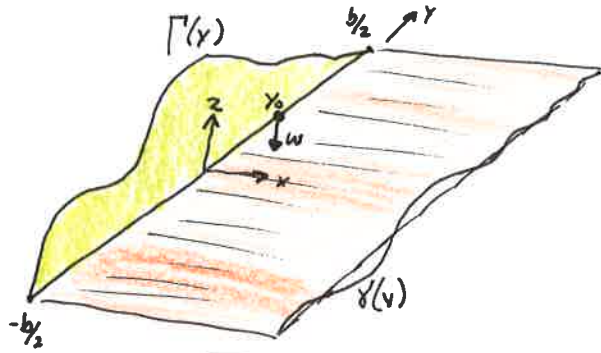
In the limit

$$\frac{dw}{dy} = -\frac{1}{4\pi} \cdot \frac{1}{y_0 - y} \cdot \frac{d\Gamma}{dy}$$

$$\Rightarrow dw = -\frac{1}{4\pi(y_0 - y)} \frac{d\Gamma}{dy} dy$$

Notice that this is only valid on the y axis ( $x=0$ )

Vertical Velocity at a single point resulting/consistent/imposed/induced by a distribution of  $\Gamma(y)$  from a wing



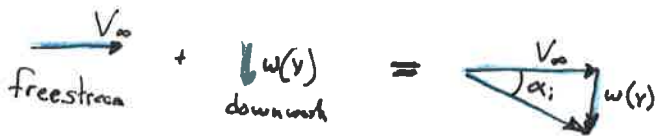
Integrate  $dw$  across the wing span

$$w(y_0) = \int dw$$

$$= \int_{-b/2}^{b/2} -\frac{1}{4\pi(y-y_0)} \frac{d\Gamma}{dy} dy$$

$$w(y_0) = -\frac{1}{4\pi} \int_{-b/2}^{b/2} \frac{d\Gamma/dy}{y-y_0} dy$$

Local angle of attack:



$\alpha_i \equiv$  Induced Angle of Attack

$$= -\arctan\left(\frac{w(y)}{V_\infty}\right) \approx -\frac{w(y)}{V_\infty}$$

What about a freestream at an angle of attack  $\alpha$ ?



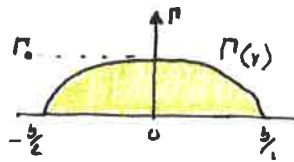
Effective angle of attack

$$\alpha_{eff} = \alpha - \alpha_i$$

$$\alpha_{eff}(y) = \alpha - \frac{w(y)}{V_\infty}$$

↑  
varies with  $y$

# Elliptical Lift Distribution



$$\Gamma(y) = \Gamma_0 \sqrt{1 - (y/b)^2} = \Gamma_0 \sqrt{1 - (2y/b)^2}$$

Lift

$$\begin{aligned} L &= \int_{-b/2}^{b/2} L' dy = \int_{-b/2}^{b/2} \rho V_\infty \Gamma_0 \sqrt{1 - (2y/b)^2} dy \quad \text{if } y = \frac{b}{2} \cos \theta \quad \text{and } dy = -\frac{b}{2} \sin \theta d\theta \\ &= \rho V_\infty \Gamma_0 \int_{\pi}^0 \sqrt{1 - \cos^2 \theta} \left(-\frac{b}{2} \sin \theta\right) d\theta = \rho V_\infty \Gamma_0 \left(-\frac{b}{2}\right) \int_{\pi}^0 \sin^2 \theta d\theta \\ &= \rho V_\infty \Gamma_0 \left(-\frac{b}{2}\right) \left. \frac{\theta}{2} - \frac{1}{2} \sin \theta \cos \theta \right|_{\pi}^0 \\ &= \frac{1}{2} \rho V_\infty \Gamma_0 \frac{b}{2} \pi \end{aligned}$$

Downwash

$$\frac{d\Gamma}{dy} = \frac{d}{dy} \left( \Gamma_0 \sqrt{1 - (2y/b)^2} \right) = -\frac{4y}{b} \frac{\Gamma_0}{\sqrt{b^2 - 4y^2}}$$

$$w(y_0) = -\frac{1}{4\pi} \int_{-b/2}^{b/2} \frac{d\Gamma/dy}{y - y_0} dy = -\frac{4}{4\pi} \frac{\Gamma_0}{b} \int_{-b/2}^{b/2} \frac{y}{\sqrt{b^2 - 4y^2}} \frac{1}{y - y_0} dy$$

Substitute  $y = \frac{b}{2} \cos \theta$

$$= -\frac{4}{4\pi} \frac{\Gamma_0}{b} \int_{\pi}^0 \frac{\frac{b}{2} \cos \theta}{\sqrt{b^2 - 4 \frac{b^2}{4} \cos^2 \theta}} \cdot \frac{1}{\frac{b}{2} \cos \theta - \frac{b}{2} \cos \theta_0} \left(-\frac{b}{2} \sin \theta\right) d\theta$$

$b \sqrt{1 - \cos^2 \theta}$

$$= \frac{\Gamma_0}{\pi b} \int_{\pi}^0 \frac{1}{2} \frac{\cos \theta}{\cos \theta - \cos \theta_0} d\theta$$

Remember this integral?

# Glauert Integral

$$\int_0^\pi \frac{\cos n\theta \, d\theta}{\cos\theta - \cos\theta_0} = \frac{\pi \sin n\theta_0}{\sin\theta_0}$$

$$w(y_0) = \frac{-\Gamma_0}{\pi b^2} \int_0^\pi \frac{\cos\theta}{\cos\theta - \cos\theta_0} d\theta = \frac{-\Gamma_0}{2\pi b} \frac{\pi \sin n\theta_0}{\sin\theta_0} = -\frac{\Gamma_0}{2b}$$

The downwash is a constant for an elliptical lift distribution

Induced angle due to lift

$$\alpha_i = -\frac{w}{V_\infty} = \frac{\Gamma_0}{2bV_\infty}$$

and  $L = \frac{1}{4} \rho V_\infty \Gamma_0 b \pi$   
 solve for  $\Gamma_0 = \frac{4L}{\rho V_\infty b \pi}$

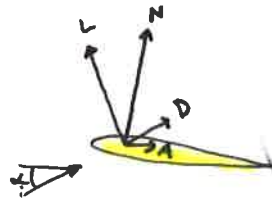
$$= \frac{4L \cdot 1}{\rho V_\infty b \pi 2bV_\infty} = \frac{L}{\frac{1}{2} \rho V_\infty^2 b^2 \pi}$$

$$C_L = \frac{L}{\frac{1}{2} \rho V_\infty^2 S} \Rightarrow C_L S = \frac{L}{\frac{1}{2} \rho V_\infty^2}$$

$$= \frac{C_L S}{\pi b^2} = \frac{C_L}{\pi AR}$$

$$\alpha_i = \frac{C_L}{\pi AR}$$

elliptical only



For small  $\alpha$ ,  $L \approx N$  since  $\cos\alpha \approx 1$   
 and  $D \approx \sin\alpha \approx \alpha$  since  $\alpha \ll 1$

Thus,  $D \approx L \alpha_i$

$$C_D \approx C_L \alpha_i = \frac{C_L^2}{\pi AR}$$

Drag due to lift, "induced drag", depends on lift squared and inversely with aspect ratio.

# Geometry of Elliptical Lift Distributions

$$\Gamma(y) = \Gamma_0 \sqrt{1 - \left(\frac{y}{b}\right)^2} \quad \text{and} \quad \Gamma_0 = \frac{4L}{\rho V_\infty b \pi}$$

$$= \frac{4L}{\rho V_\infty b \pi} \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

Also,

$$L' = C_L \cdot g \cdot c(y) \Rightarrow c(y) = \frac{L'}{C_L g}$$

$\uparrow$  lift per span       $\uparrow$  lift coeff       $\uparrow$  dynamic pressure      chord

Since  $\alpha_i$  is constant for an elliptical wing,  $C_L$  is only  $C_L(\alpha)$  (not a function of  $y$ )  
 $C_L = 2\pi(\alpha - \alpha_i - \alpha_{2c})$

Also

$$L' = \rho V \Gamma(y) = \rho V \frac{4L}{\rho V b \pi} \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

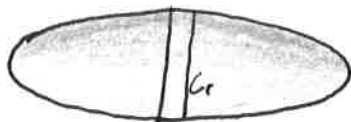
$$= \frac{4L}{b \pi} \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

Thus

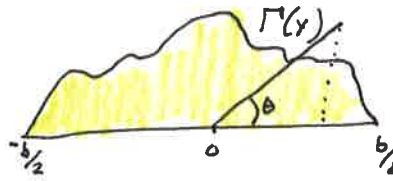
$$c(y) = \frac{L'}{C_L g} = \frac{4L}{b \pi} \sqrt{1 - \left(\frac{2y}{b}\right)^2} \cdot \frac{1}{2\pi(\alpha - \alpha_i - \alpha_{2c})} \cdot \frac{1}{\frac{1}{2} \rho V_\infty^2}$$

The only  $y$  variation is  $\sqrt{1 - \left(\frac{2y}{b}\right)^2}$

$$c(y) = \sqrt{1 - \left(\frac{2y}{b}\right)^2} \cdot C_D \quad \text{This describes an ellipse.}$$



# General Lift Distribution



Expand  $\Gamma(y)$  in terms of Fourier components

$$y = \frac{b}{2} \cos \theta \quad dy = -\frac{b}{2} \sin \theta d\theta$$

$$\Gamma(\theta) = 2bV_\infty \sum_{n=1}^{\infty} A_n \sin(n\theta)$$

Notice that we picked  $\sin n\theta$  which always fits the boundary conditions of  $\Gamma(0) = \Gamma(\pi) = 0$

$$A_1 \sin \theta + A_2 \sin 2\theta + A_3 \sin 3\theta + A_4 \sin 4\theta$$

Downwash eqn needs  $d\Gamma/dy$

$$\frac{d\Gamma}{dy} dy = \left( \frac{d\Gamma}{d\theta} \frac{d\theta}{dy} \right) \left( d\theta \frac{dy}{d\theta} \right) = \frac{d\Gamma}{d\theta} d\theta$$

$$= 2bV_\infty \sum_{n=1}^{\infty} (A_n \cos(n\theta) n) d\theta$$

Wake velocity

$$w(y_0) = -\frac{1}{4\pi} \int_{-b/2}^{b/2} \frac{d\Gamma/dy}{y-y_0} dy = -\frac{1}{4\pi} \int_{\pi}^0 \frac{2bV_\infty \sum (A_n \cos n\theta n) d\theta}{\frac{b}{2} \cos \theta - \frac{b}{2} \cos \theta_0}$$

Did you catch the trick?  
Amazing!

$$= \frac{V_\infty}{\pi} \int_0^\pi A_n \frac{\cos n\theta}{\cos \theta - \cos \theta_0} d\theta = \frac{V_\infty}{\pi} \sum_{i=1}^{\infty} n A_n \int_0^\pi \frac{\cos n\theta}{\cos \theta - \cos \theta_0} d\theta$$

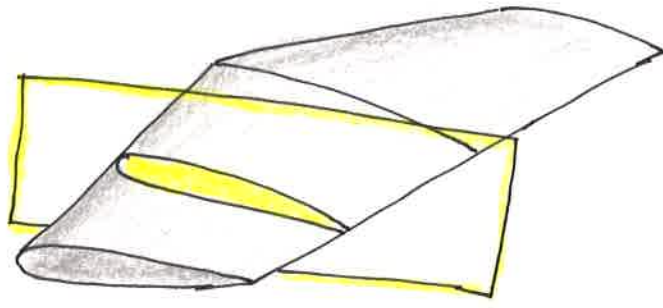
Glauert Integral

$$\frac{\pi \sin n\theta_0}{\sin \theta_0}$$

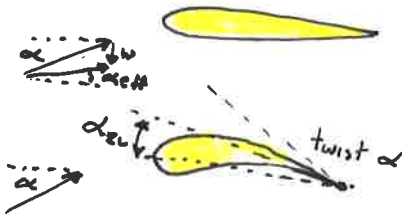
$$= \frac{V_\infty}{\pi} \pi \sum_{i=1}^{\infty} n A_n \frac{\sin n\theta_0}{\sin \theta_0}$$

$$w(\theta_0) = V_\infty \sum_{i=1}^{\infty} n A_n \frac{\sin n\theta_0}{\sin \theta_0}$$

$$\alpha_i = \frac{w}{V_\infty}$$



On this cross section,



$$\text{Effective AOA} = \alpha + \alpha_{\text{aero}} - \alpha_i$$

$\alpha_{\text{aero}} = \alpha_{\text{geo}} + \alpha_{2c}$

$$L' = \rho V \Gamma$$

$$a \frac{L'}{\rho c} = \frac{\rho V \Gamma}{\frac{1}{2} \rho V^2 c}$$

$$C_L = \frac{2 \Gamma}{V c}$$

$$\Gamma = \frac{1}{2} V c C_L$$

Expand  $\Gamma$  w Fourier series

$$\Gamma = 2b V_{\infty} \sum_{n=1}^{\infty} A_n \sin(n\theta) = \frac{1}{2} V_{\infty} c(\theta) C_L(\alpha, \alpha_{\text{aero}}, \alpha_i)$$

Circulation at a particular point  $\theta$  in the wing      chord function wrt  $\theta$       lift coefficient

Linear Aerodynamics

$$C_L = 2\pi(\alpha + \alpha_{\text{aero}} - \alpha_i)$$

Expand

$$2b V_{\infty} \sum_{n=1}^{\infty} A_n \sin(n\theta) = \frac{1}{2} V_{\infty} c(\theta) C_{L\alpha} \left( \alpha + \alpha_{\text{aero}}(\theta) - \sum_{n=1}^{\infty} \left( n A_n \frac{\sin n\theta}{\sin \theta} \right) \right)$$

Rearrange

$$\sum_{n=1}^{\infty} A_n \left( \sin(n\theta) + \frac{c(\theta)}{4b} C_{L\alpha} n \frac{\sin n\theta}{\sin \theta} \right) = \frac{c(\theta)}{4b} C_{L\alpha} (\alpha + \alpha_{\text{aero}})$$

freestream + downwash is tangent to the local effective airfoil section

Solve for  $A_n$  given  $c(\theta)$ ,  $\alpha_{\text{aero}}(\theta)$ ,  $\alpha$ ,  $b$



Lift

$$L = \int L' dy = \int p V \Gamma dy = p V_\infty \int_{-\pi}^{\pi} 2b V_\infty \sum (A_n \sin n\theta) \left(-\frac{b}{2} \sin\theta d\theta\right)$$

$$= \frac{2b^2 V_\infty^2 p}{2} \int_{-\pi}^{\pi} - \underbrace{\left(\sum A_n \sin n\theta\right)}_{\text{orthogonal, so only } n=1 \text{ contributes!!}} \sin\theta d\theta$$

$$= b^2 V_\infty^2 p \frac{\pi}{2} A_1$$

$$C_L = \frac{L}{\frac{1}{2} \rho V^2 S} = \frac{b^2 V_\infty^2 p \frac{\pi}{2} A_1}{\frac{1}{2} \rho V^2 S} = \pi A R A_1$$

Lift only depends on the  $A_1$  term.

What is the  $A_1$  term?

$$\Gamma = 2b V_\infty A_1 \sin\theta$$

convert to  $y$  coordinate

$$\Gamma = 2b V_\infty A_1 \sin\left(\arccos\left(\frac{2y}{b}\right)\right)$$

$$y = \frac{b}{2} \cos\theta$$

$$\frac{2y}{b} = \cos\theta$$

$$\theta = \arccos\left(\frac{2y}{b}\right)$$

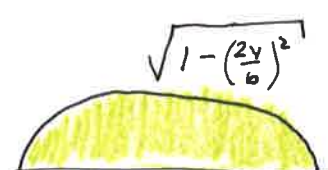
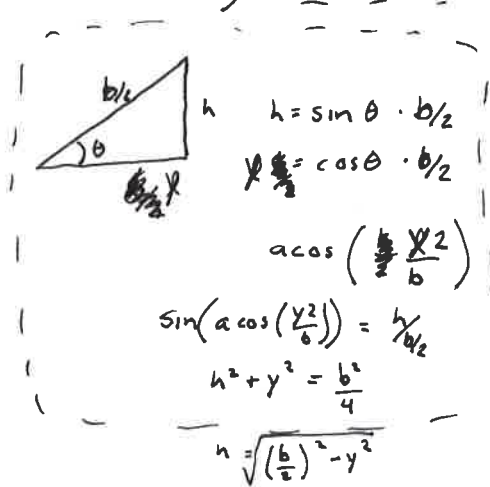
$$\Gamma = 2b V_\infty A_1 \frac{\sin\left(\arccos\left(\frac{2y}{b}\right)\right)}{\sqrt{\left(\frac{b}{2}\right)^2 - y^2}}$$

$$= 2b V_\infty A_1 \sqrt{\left(\frac{b}{2}\right)^2 - y^2} \cdot \frac{2}{b} \cdot \frac{b}{2}$$

$$= 2b V_\infty A_1 \sqrt{1 - \left(\frac{2y}{b}\right)^2} \cdot \frac{b}{2}$$

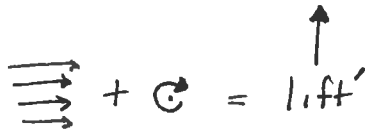
$$= b^2 V_\infty A_1 \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

Elliptical Distribution

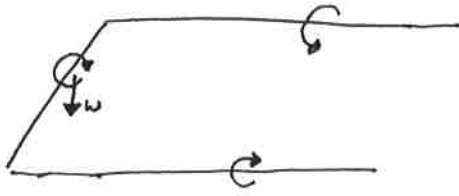


The overall lift is the elliptical portion of  $\Gamma$

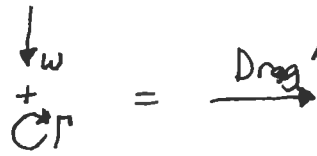
# Drag



What if we consider drag' as the force in the perpendicular direction from the vertical perturbation / downwash velocity?



$$\text{Drag}' \approx w \Gamma$$



Integrate  $D'$  over span.

$$D = \int D' dy = \int_{-b/2}^{b/2} w \Gamma dy = \int_{-\pi}^0 V_{\infty} \left( \sum_{i=1}^{\infty} n A_n \frac{\sin n \theta}{\sin \theta} \right) \left( 2b V_{\infty} \sum_{i=1}^{\infty} A_n \sin n \theta \right)^{-1/2} \sin \theta d\theta$$

$$= V_{\infty}^2 b^2 \int_0^{\pi} \underbrace{\left( \sum_{i=1}^{\infty} n A_n \frac{\sin n \theta}{\sin \theta} \right) \left( \sum_{i=1}^{\infty} A_n \sin n \theta \right)}_{\text{term by term orthogonal}} \sin \theta d\theta$$

$$D = \frac{1}{2} \rho V^2 b^2 \pi \sum_{i=1}^{\infty} n A_n^2$$

$$C_D = \frac{D}{\frac{1}{2} \rho V^2 S} = \frac{\frac{1}{2} \rho V^2 b^2 \pi \sum n A_n^2}{\frac{1}{2} \rho V^2 S} = AR \pi \sum n A_n^2$$

$$C_L = \pi AR A_1$$

$$C_D = AR \pi \left[ \sum n \left( \frac{A_n}{A_1} \right)^2 \right] \cdot A_1^2 = AR \pi \left[ \sum n \left( \frac{A_n}{A_1} \right)^2 \right] \frac{C_L^2}{\pi^2 AR^2}$$

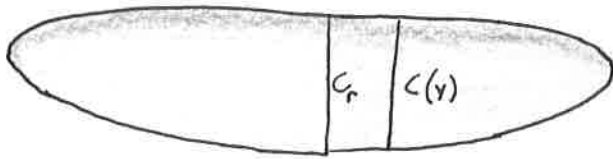
Drag depends on  $C_L^2$  and inversely with  $AR$  and with squares of non-elliptical terms

$$C_D = \frac{1}{\pi AR} C_L^2 (1 + \delta)$$

$$\delta = \sum_{n=2}^{\infty} n \left( \frac{A_n}{A_1} \right)^2$$

$$1 + \delta \equiv \frac{1}{e}$$

# Flat Elliptical Wing



$$A_1 = \frac{\frac{C_{L\alpha}}{\pi AR}}{1 + \frac{C_{L\alpha}}{\pi AR}} \cdot (\alpha + \alpha_{zero})$$

- Elliptical
- No twist "flat"

$$c(\theta) = C_r \sin \theta$$

$$c(y) = C_r \sqrt{1 - \left(\frac{2y}{b}\right)^2}$$

$$S = \frac{\pi}{4} b C_r$$

$$AR = \frac{4}{\pi} \frac{b}{C_r}$$

$$C_{L\alpha} = \frac{C_{L\alpha}}{1 + \frac{C_{L\alpha}}{\pi AR}}$$

$$C_{Di} = \frac{C_L^2}{\pi AR}$$

# Oswald Efficiency Factor

$$C_D = \frac{1}{\pi AR} \cdot \frac{1}{e} \cdot C_L^2$$

↑ Oswald Eff' factor  $\equiv e \equiv \frac{1}{1+\delta}$

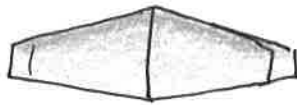
"e" measures the efficiency with respect to <sup>induced</sup> drag as compared to an elliptical wing



This is a powerful approach valid for more than lifting line assumptions

## Straight wings

From: Aircraft Design, Raymer



$$e \approx 1.78 (1 - 0.045 A^{0.68}) - 0.64$$

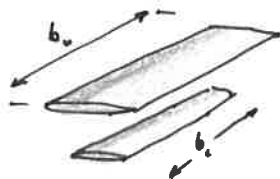
## Swept wings



$$e \approx 4.61 (1 - 0.045 A^{0.68}) (\cos \Lambda_{LE})^{0.15} - 3.1$$

$\Lambda_{LE} \geq 30^\circ$

## Biplanes



$$e \approx \frac{\mu^2 (1+r^2)}{\mu^2 + 2\sigma\mu r + r^2}$$

$$\mu = \frac{\text{shorter span}}{\text{longer span}}$$

$\sigma$  = Interference factor depending on gap / span

$$r = \frac{\text{lift on shorter}}{\text{lift on longer}}$$

$$AR = \frac{\max(b_u, b_l)}{S_u + S_l}$$

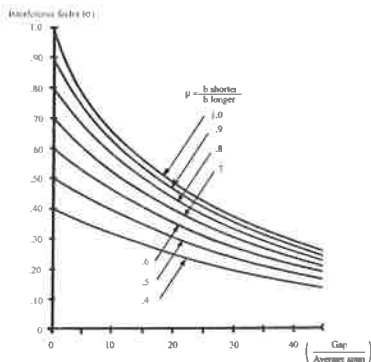


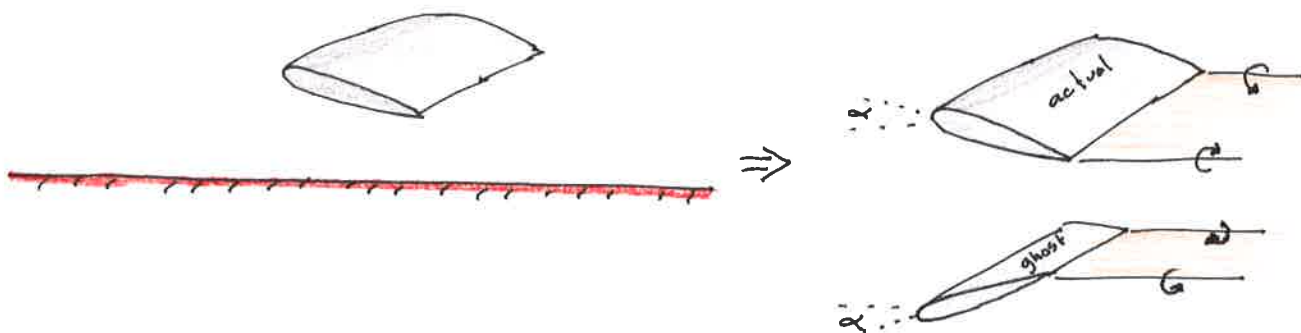
Fig. 12.34 Prandtl's biplane interference factor (Ref. 12).

## Ground Effect:

Near the ground, a wing becomes more efficient.

$$C_D = K_{eff} C_L^2 \quad \text{with} \quad \frac{K_{eff}}{K} = \frac{33 \left(\frac{h}{b}\right)^{1.5}}{1 + 33 \left(\frac{h}{b}\right)^{1.5}}$$

We can model this with a ghost image of a wing below the surface

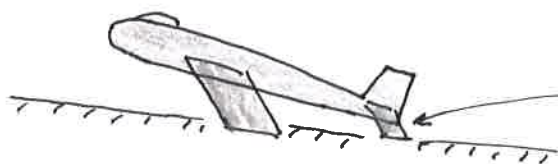


- In effect (pun intended), the downwash is prevented from passing through the ground
- Or, consider the ghost image has an upwash on the actual wing

We saw this in the [tiny.cc/AEM313](https://www.tiny.cc/AEM313) Wings Vortex video, where the vortex moved outboard ~~again~~ as the aircraft approached the ground.

Ground effect can be a serious issue for the unsuspecting pilot (especially in low wing Piper a/c)

Ground effect can affect S+C at takeoff and landing



The tail is nearer the ground as the aircraft rotates for takeoff

Ground effect allows for high speed flying boats (Ground Effect Vehicle)

"Ekranoplan"

[tiny.cc/AEM313](https://www.tiny.cc/AEM313) Ground Effect