

L15p1

$$U_t = U_{xx} - 2U_x$$

$$U(x,0) = \sin(x)$$



Visually identify diffusion terms  $U_{xx}$  and convection terms  $-2U_x$

Strategy: Transform from  $x$  coordinate to a new  $\xi$  coordinate traveling at velocity 2.

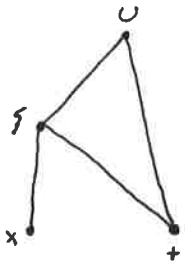
$$\xi = x - 2t$$

$\frac{du}{dt}$  doesn't change since only  $x$  is transformed, right? Charge ahead!

$$\frac{du}{dx} = \frac{du}{d\xi} \frac{d\xi}{dx} = \frac{du}{d\xi} \cdot 1 = \frac{du}{d\xi} \Rightarrow \frac{d^2u}{dx^2} = \frac{d^2u}{d\xi^2}$$

PDE:  $\frac{du}{dt} = U_{\xi\xi} - 2U_{\xi}$  Nothing happened! What is wrong?

Try again, be consistent with transform.



$$\frac{du}{dx} = \frac{du}{d\xi} \frac{d\xi}{dx} \quad \text{and} \quad \frac{d^2u}{dx^2} = \frac{d^2u}{d\xi^2}$$

$$\frac{du}{dt} = \frac{du}{d\xi} \frac{d\xi}{dt} + \frac{du}{dt} = \frac{du}{d\xi} (-2) + \frac{du}{dt}$$

Notice this term!

Subs into PDE

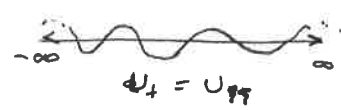
$$\frac{du}{dt} - 2 \frac{du}{d\xi} = \frac{d^2u}{d\xi^2} - 2U_x \Rightarrow \frac{du}{dt} = \frac{d^2u}{d\xi^2}$$

(Note: The term  $-2U_x$  is crossed out with a bracket and a zero below it, indicating it cancels out.)

So in the  $\xi$  direction/coordinate, this is only a diffusion problem.

Use Fourier technique for  $\xi = -\infty \Leftrightarrow \infty$

$$\frac{d}{dt}(F(\omega) \cdot U) = -\omega^2 U \Rightarrow U(t) = e^{-\omega^2 t} U_0$$



and  $F(\text{IC}) = F(\sin(x)) = \underset{\text{back at book}}{\text{back at book}} = i\sqrt{\frac{\pi}{2}} (\delta(\omega+1) - \delta(\omega-1))$

$$U = \overset{\text{DDE}}{e^{-\omega^2 t}} \overset{\text{FC}}{i\sqrt{\frac{\pi}{2}} (\delta(\omega+1) - \delta(\omega-1))}$$

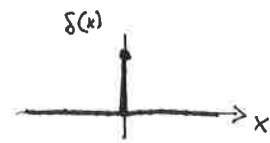
$U = F^{-1}(U)$  but  $\uparrow$  is certainly not in the book! So messy, right?!

Look at the definition of  $F^{-1}(f) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega$

This would look like a nasty integral. But we have  $\delta$  functions!!

$$U = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-\omega^2 t} i\sqrt{\frac{\pi}{2}} (\delta(\omega+1) - \delta(\omega-1)) e^{i\omega x} d\omega$$

only has value when  $\omega = -1$ 
only has value when  $\omega = +1$



$$= \underbrace{\frac{i}{2} e^{-t} e^{-ix}}_{\omega = -1} - \underbrace{\frac{i}{2} e^{-t} e^{ix}}_{\omega = +1} = \frac{i}{2} e^{-t} (e^{-ix} - e^{ix}) \cdot \frac{i}{1} = \frac{e^{-t}}{2i} (e^{ix} - e^{-ix})$$

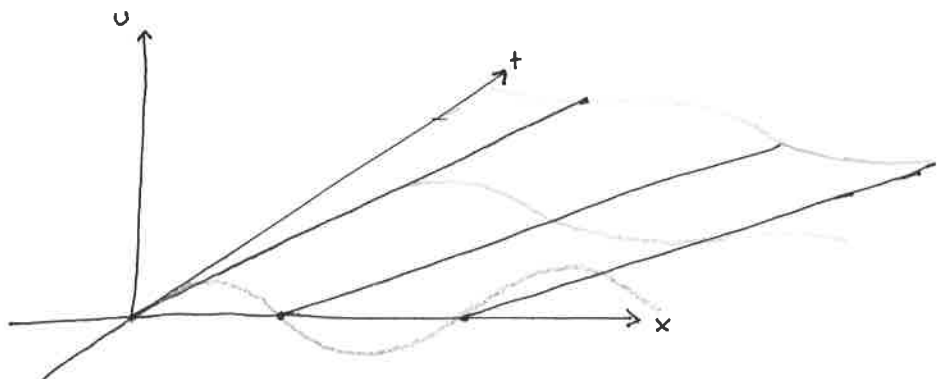
Recognize this?!

$$= e^{-t} \sin(x) \quad \leftarrow \text{technically, this is } \xi \quad = e^{-t} \sin(\xi)$$

Transform back to  $x$  coordinates  $x = \xi + 2t$  or  $\xi = x - 2t$

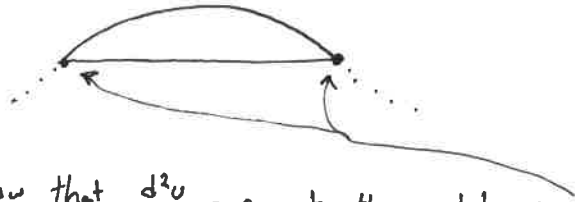
$$U = e^{-t} \sin(x - 2t)$$

Translating sine wave that decreases in magnitude in time.



Faster solution

Isolate one cell of sine wave



We know that  $\frac{d^2 u}{d \xi^2} = 0$  at the nodal points

Thus, the solution should be the finite domain solution as well.

$$U\left(\frac{\xi}{\pi}, 0\right) = \sin\left(\frac{\xi}{\pi}\right) \quad 0 < \xi < \pi$$

Cell solution

$$U(x, t) = e^{-t} \sin\left(\frac{x}{\pi}\right) \quad \text{in } 0 < x < \pi$$

Tile solution ( $\pm$ ) across domain.

$$U = e^{-t} \sin\left(\frac{x}{\pi}\right)$$

Transform  $\xi \rightarrow x - 2t$

$$U = e^{-t} \sin(x - 2t)$$

15p2

$$U_t = U_{xx} - 2U_x$$

$$U(x,0) = e^x \sin(x)$$

$U_t = U_{xx}$  diffusion and  $U_t = -2U_x$  convection

Solution  $\approx U = e^{k_1(x-ct)}$

pull this out of  $U(x,t)$

$$U(x,t) = e^{k_1(x-ct)} w(x,t)$$

Subst. into gov. eqn.

$$\underbrace{-k_1 c e^{k_1(x-ct)} w + e^{k_1(x-ct)} w_t}_{U_t} = \underbrace{k_1^2 e^{k_1(x-ct)} w + 2k_1 e^{k_1(x-ct)} w_x + e^{k_1(x-ct)} w_{xx}}_{U_{xx}}$$

pull out  $e^{k_1(x-ct)}$  term  $-c = \frac{1}{2} \Rightarrow c=1$

$$\underbrace{-k_1 c w + w_t}_{U_t} = \underbrace{k_1^2 w + 2k_1 w_x + w_{xx} - 2k_1 w - 2w_x}_{U_{xx}}$$

$$\underbrace{-k_1 c w + w_t}_{U_t} = \underbrace{k_1^2 w + 2k_1 w_x + w_{xx}}_{U_{xx}} \underbrace{- 2k_1 w - 2w_x}_{-2U_x}$$

if  $k_1 = 1$

If  $c=1$  and  $k_1=1$

$$\underline{w_t = w_{xx}}$$

IC

$$U(x,0) = e^x \sin(x) = e^{k_1(x-ct)} w(x,0) = e^x w(x,0) \Rightarrow w(x,0) = \sin(x)$$

Solution to  $w$

~~$w = e^{-t} \sin(x)$~~ 

$$w = e^{-t} \sin(x)$$

Combine back to  $u = e^{(x-t)} w$

$$u = e^{(x-t)} e^{-t} \sin(x)$$

$$u = e^x e^{-2t} \sin(x)$$

L13p3  $\bar{u}$  Duhamel

$$U_t = U_{xx} \quad 0 < x < \infty$$

$$U(0,t) = \sin(t)$$

$$U(x,0) = 0$$

$$L(U_t) = L(U_{xx})$$

$$sU - \cancel{U(x,0)}^0 = U_{xx} \quad \Rightarrow \quad U = Ae^{-\sqrt{s}x}$$

Duhamel's method says to find the impulse response.  $U(0,t) = \delta(t)$  so  $f(t) = \sin(t)$

$$L(U(0,t)) = L(\delta(t)) = |e^{-s \cdot 0}| = 1$$

$$f'(t) = \cos(t)$$

$$U = e^{-\sqrt{s}x}$$

$$\mathcal{L}^{-1}(U) = \mathcal{L}^{-1}(e^{-\sqrt{s}x}) = \frac{x}{2\sqrt{\pi t}} e^{-\frac{x^2}{4t}} \quad \text{Not in the book...}$$

$$w = \frac{x}{2\sqrt{\pi t}} e^{-\frac{x^2}{4t}}$$

Duhamel:  $U(x,t) = \int_0^t w_+(x,t-\tau) f(\tau) d\tau$  ← Not as useful as below ( $w_+$  is difficult)

and

$$= \int_0^t w(x,t-\tau) f'(\tau) d\tau + f(0)w(x,t)$$

← we want to use this form

$$= \int_0^t \frac{x}{2\sqrt{\pi(t-\tau)}} e^{-\frac{x^2}{4(t-\tau)}} \cos(\tau) d\tau + \cancel{f(0)}^0 w(x,t)$$

$$U(x,t) = \int_0^t \frac{x}{2\sqrt{\pi(t-\tau)}} e^{-\frac{x^2}{4(t-\tau)}} \cos(\tau) d\tau$$

L6 p3

$$U_t = U_{xx} \quad 0 < x < 1$$

$$U_x(0,t) = 0$$

$$U_x(1,t) + h U(1,t) = 1$$

$$0 < t < \infty$$

$$U(x,0) = \sin(\pi x)$$

The  $x=1$  BC is not homogeneous.

- Break into S.S and transient components

$$U = \bar{U} + \mathcal{U} \quad \text{with } \bar{U} = Ax + B$$

1) PDE  $\bar{U}_t + \mathcal{U}_t = \bar{U}_{xx} + \mathcal{U}_{xx} \Rightarrow \boxed{\mathcal{U}_t = \mathcal{U}_{xx}}$

2) BCs  $\bar{U}_x(0,t) + \mathcal{U}_x(0,t) = 0 \Rightarrow A + \mathcal{U}_x(0,t) = 0 \quad A = 0$   
 $\boxed{\mathcal{U}(0,t) = 0}$

$$\bar{U}_x(1,t) + \mathcal{U}_x(1,t) + h \bar{U}(1,t) + h \mathcal{U}(1,t) = 1$$

$$A + \mathcal{U}_x(1,t) + h(0 \cdot 1 + B) + h \mathcal{U}(1,t) = 1$$

$$\mathcal{U}_x(1,t) + h \mathcal{U}(1,t) = 1 - hB \Rightarrow B = \frac{1}{h}$$

$$\boxed{\mathcal{U}_x(1,t) + h \mathcal{U}(1,t) = 0}$$

with  $\boxed{\bar{U} = \frac{1}{h}}$

- 3) IC

$$\bar{U} + \mathcal{U} = \sin(\pi x)$$

$$\Rightarrow$$

$$\boxed{\mathcal{U} = \sin(\pi x) - \frac{1}{h}}$$

L6P3 continued

- Solve new PDE

$$U = \sum X T \quad \text{with} \quad \begin{aligned} U_t &= U_{xx} \\ U(0,t) &= 0 \\ U_x(1,t) + hU(1,t) &= 0 \\ U &= \sin(\pi x) - \frac{1}{h} \end{aligned}$$

- Sep of Vars'

$$X T_t = X_{xx} T \quad \Rightarrow \quad \frac{T_t}{T} = \frac{X_{xx}}{X} = -\lambda^2$$

$$\text{1st order } T \Rightarrow T = T_0 e^{-\lambda^2 t}$$

$$\text{2nd order } X \Rightarrow X = A_m \sin \lambda_m x + B_m \cos \lambda_m x$$

- Apply BCs

$$U(0,t) = 0 = X(0)T(t) \quad \Rightarrow \quad X(0) = 0 = A_m \sin 0 + B_m \cos 0$$

Thus  $B_m = 0$

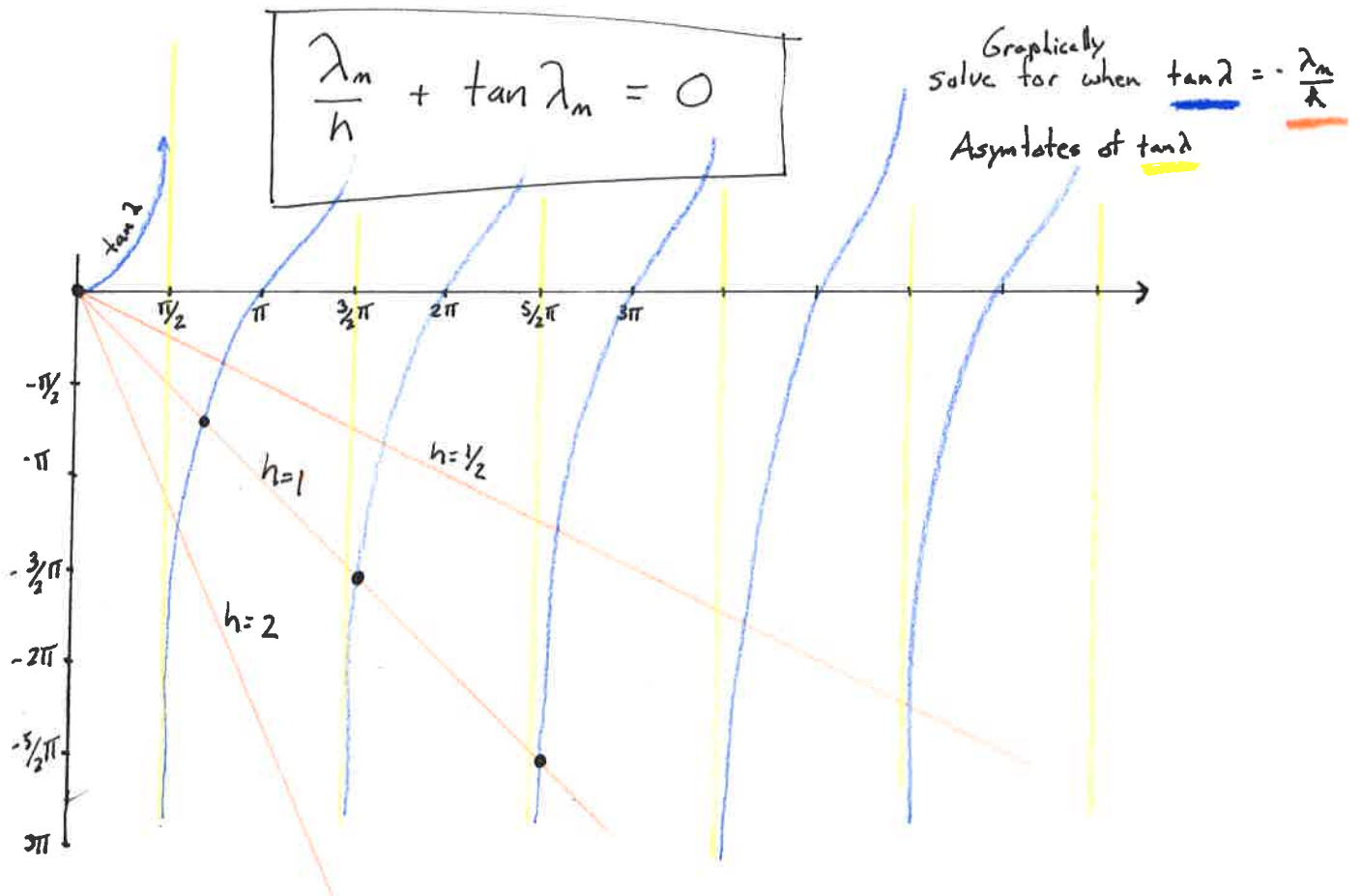
$$\begin{aligned} U_x(1,t) + hU(1,t) &= 0 = X_x(1,t) + hX(1,t) \\ &= \lambda A_m \cos \lambda_m x + h A_m \sin \lambda_m x \\ &= A_m (\lambda_m \cos \lambda_m + h \sin \lambda_m) \end{aligned}$$

So, either  $A_m = 0$  or  $\lambda_m \cos \lambda_m + h \sin \lambda_m = 0$

trivial  
boring  
solution

- Find  $\lambda_n$

$$\lambda_m \cos \lambda_m + h \sin \lambda_m = 0$$



- Find  $\lambda_n$  when  $h=1$

$$\lambda_0 = 0 \text{ useless since } \sin 0 = 0$$

$$\lambda_1 \approx \text{halfway between } \frac{\pi}{2} \text{ and } \pi \approx \frac{3}{4}\pi \approx 2.3$$

$$\lambda_2 \approx \text{asymptote at } \frac{3}{2}\pi \approx 4.7$$

$$\lambda_m \approx m_{th} \text{ asymptote of } \tan \lambda \approx \left(\frac{2m-1}{2}\right)\pi$$

Exact

0

2.02876

4.91318



• Solution

$$X = A_m \sin(\lambda_m x)$$

$$= A_0 \sin 0 + A_1 \sin(2.3x) + A_2 \sin(4.7x) + \dots + A_n \sin\left(\frac{2n-1}{2}\pi x\right)$$

• Find  $A_n$  terms.

$$\phi = A_m \sin(\lambda_m x)$$

premultiply by  $\sin(\lambda_n x)$

Integrate over domain  $\int_0^1 \dots dx$

$$\int_0^1 \phi(x) \sin(\lambda_n x) dx = A_m \int_0^1 \sin(\lambda_m x) \sin(\lambda_n x) dx$$

$$A_m = \frac{\int_0^1 \left(\sin \pi x - \frac{1}{n}\right) \sin(\lambda_n x) dx}{\int_0^1 \sin(\lambda_n x) \sin(\lambda_n x) dx}$$

↖ From S-L theory, this only is non-zero when  $n=m$ !

• Solution

$$U = \bar{U} + \psi = \frac{1}{h} + A_1 e^{-4.113t} \sin(2.02x) + A_2 e^{-24.14t} \sin(4.7x) + \dots$$

$$\text{where } A_m = \frac{\int_0^1 \left(\sin(\pi x) - \frac{1}{h}\right) \sin(\lambda_m x) dx}{\int_0^1 \sin^2(\lambda_m x) dx}$$

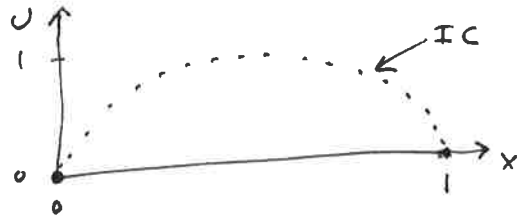
L9 p1

$$U_t = U_{xx} + \sin(3\pi x)$$

$$U(0, t) = 0$$

$$U(1, t) = 0$$

$$U(x, 0) = \sin(\pi x)$$



This is a non-homogeneous PDE ( $\sin 3\pi x$ ). Call  $f_n = \sin 3\pi x$

• Find eigenfunctions by solving homogeneous PDE

$$\begin{aligned} U_t &= U_{xx} \\ U(0) &= 0 \\ U(1) &= 0 \end{aligned}$$

by inspection  $\Rightarrow$

$$U = X T$$

with  $X = \sin(n\pi x)$   
these are our building blocks...

• Expand ~~non-homogeneous~~ non-homogeneous PDE with

$$U = T_n(t) X(x)$$

$$T_{n,t} X + T_n X_t = T_{n,xx} X + T_n X_{xx} + f_n$$

• Expand  $f_n$  in terms of  $X$ :  $f_n = \sin(n\pi x)$

$$\int_0^1 \underset{\substack{\uparrow \\ \sin(3\pi x)}}{f_n} \sin(n\pi x) dx = \int_0^1 \sin(m\pi x) \sin(n\pi x) dx$$

Obvious?

$$f_n = \begin{cases} 1 & n=3 \\ 0 & \text{otherwise} \end{cases}$$

$$f_n = 0 \cdot \sin(\pi x) + 0 \cdot \sin(2\pi x) + 1 \cdot \sin(3\pi x) + \dots$$

• Continue with PDE

$$T_{n,t} X + \cancel{T_n X_t} = \cancel{T_n X_{xx}} + T_n \cancel{X_{xx}} + \left. \begin{array}{l} -n^2 \pi^2 \sin(n\pi x) = -n^2 \pi^2 X \\ 1 \quad n=3 \\ 0 \quad \text{otherwise} \end{array} \right\} X$$

• BCs

$$\begin{aligned} U(0,t) = 0 &= T_n(t) \cancel{X(0)} \\ U(1,t) = 0 &= T_n(t) \cancel{X(1)} \end{aligned} \Rightarrow \text{Not useful. Why? } X!$$

• ICs

$$U(x,0) = \sin(\pi x) = T_n(0) X(x)$$

In general, premultiply by  $X$  and integrate to find  $T_n(0)$

$$\int_0^1 X \sin(\pi x) dx = \int_0^1 T_n(0) X^2(x) dx \quad \text{with } X(x) = \sin(n\pi x)$$

$$T_n(0) = \begin{cases} 1 & n=1 \\ 0 & \text{otherwise} \end{cases}$$

• Solve for  $T_n(t)$  in PDE

$$X \left[ T_{n,t} = -n^2 \pi^2 T_n(t) + \left. \begin{array}{l} 1 \quad n=3 \\ 0 \quad \text{otherwise} \end{array} \right\} \right]$$

Since  $X$  is not always zero, the bracket terms must be.

$$T_{n,t} + n^2 \pi^2 T_n = \begin{cases} 1 & n=3 \\ 0 & \text{otherwise} \end{cases}$$

$$T_n(0) = \begin{cases} 1 & n=1 \\ 0 & \text{otherwise} \end{cases}$$

$$n=1$$

$$T_{n,t} + n^2 \pi^2 T_n = 0$$

$$T_n(0) = 1$$

$$\Rightarrow T_n = T_n(0) e^{-n^2 \pi^2 t} = e^{-n^2 \pi^2 t}$$

$$n=2$$

$$T_{n,t} + n^2 \pi^2 T_n = 0$$

$$T_n(0) = 0$$

$$\Rightarrow T_n = T_n(0) e^{-n^2 \pi^2 t} = 0$$

$$n=3$$

$$T_{n,t} + n^2 \pi^2 T_n = 1$$

$$T_n(0) = 0$$

premultiply by  $e^{n^2 \pi^2 t}$

$$\Rightarrow \underbrace{e^{n^2 \pi^2 t} T_{n,t} + e^{n^2 \pi^2 t} T_n}_{\substack{\text{product rule} \\ \frac{d}{dt}(e^{n^2 \pi^2 t} T_n)}} = e^{n^2 \pi^2 t}$$

integrate wot time

$$\int_0^t \frac{d}{dt} (e^{n^2 \pi^2 t} T_n) dt = \int_0^t e^{n^2 \pi^2 t} dt$$

$$e^{n^2 \pi^2 t} T_n = \frac{e^{n^2 \pi^2 t}}{n^2 \pi^2} \Big|_0^t = \frac{e^{n^2 \pi^2 t} - 1}{n^2 \pi^2}$$

$$T_n = \frac{1 - e^{-n^2 \pi^2 t}}{n^2 \pi^2}$$

$$n=4 \quad T_n = 0$$

• Solution

$$U = e^{-\pi^2 t} \sin(\pi x) + 0 \cdot \sin(2\pi x) + \frac{1 - e^{-9\pi^2 t}}{9\pi^2} \sin(3\pi x) + 0 \dots$$

$$U = e^{-\pi^2 t} \sin(\pi x) + \frac{1 - e^{-9\pi^2 t}}{9\pi^2} \sin(3\pi x)$$

L13p2 "Tricky"

$$U_t = \alpha^2 U_{xx} \quad -\infty < x < \infty$$

$$U(x, 0) = \sin x$$

$$L(U_t) = L(\alpha^2 U_{xx}) \Rightarrow sU - \cancel{U(x, 0)}^{\sin x} = \alpha^2 U_{xx} \Rightarrow U_{xx} - \frac{s}{\alpha^2} U = -\frac{\sin x}{\alpha^2}$$

We know that  $U$  is composed of a homogeneous solution and a particular solution

$$U = A e^{-\frac{\sqrt{s}}{\alpha} x} + B e^{\frac{\sqrt{s}}{\alpha} x} + C \cdot p(x)$$

must be bounded

find a  $p(x)$ ...

expect  $p(x)$  to look like  $\sin(x) \cdot (\text{terms}) \Rightarrow p(x) = \sin(x) Z(x)$

substitute  $U = A e^{-\frac{\sqrt{s}}{\alpha} x} + C p(x)$  into ODE

$$A \frac{s}{\alpha^2} e^{-\frac{\sqrt{s}}{\alpha} x} + C (-Z) \sin(x) - \frac{s}{\alpha^2} \left( A e^{-\frac{\sqrt{s}}{\alpha} x} + C Z \sin(x) \right) = -\frac{\sin(x)}{\alpha^2}$$

0 always cancel...

Remaining terms.

$$-Z C \sin(x) \left( 1 + \frac{s}{\alpha^2} \right) = -\frac{\sin(x)}{\alpha^2}$$

Since  $\sin(x)$  is not always zero...

$$ZC = \frac{1}{\alpha^2} \frac{1}{1 + \frac{s}{\alpha^2}} = \frac{1}{\alpha^2 + s}$$

Solution to  $U$

$$U = A e^{-\frac{\sqrt{s}}{\alpha} x} + \frac{1}{\alpha^2 + s} \sin(x)$$

Inverse Laplace transform is linear

$$\mathcal{L}^{-1}(A+B) = \mathcal{L}^{-1}(A) + \mathcal{L}^{-1}(B)$$

$$\text{and } \mathcal{L}^{-1}\left(A e^{-\frac{\sqrt{s}}{\alpha} x}\right) = \mathcal{L}^{-1}\left(s \cdot \frac{e^{-\frac{\sqrt{s}}{\alpha} x}}{s}\right) = \text{complicated}$$

$$\mathcal{L}^{-1}\left(\frac{\sin(x)}{\alpha^2 + s}\right) = \sin(x) \mathcal{L}^{-1}\left(\frac{1}{\alpha^2 + s}\right) = \sin(x) e^{-\alpha^2 t}$$

Fit the IC.  $U(x, 0) = \sin(x)$

$$U = A \cdot \text{complicated} + \sin(x) e^{-\alpha^2 t}$$

only this term fits ICs  $\Rightarrow$

$$U = \sin(x) e^{-\alpha^2 t}$$

L13p3

$$U_t = U_{xx} \quad 0 < x < \infty$$

$$U(0,t) = \sin(t) \quad 0 < t < \infty$$

$$U(x,0) = 0 \quad 0 \leq x < \infty$$



Apply Laplace transform to  $t$  (BC has time term)

$$\mathcal{L}(U_t) = \mathcal{L}(U_{xx})$$

$$sU(x) - U(x,0) = \frac{d^2 U(x)}{dx^2}$$

Apply IC and simplify to canonical form

$$sU(x) = \frac{d^2 U}{dx^2} \Rightarrow U_{xx} - sU = 0$$

Apply  $\mathcal{L}$  to BC

$$\mathcal{L}(U(0,t)) = \mathcal{L}(\sin(t)) = \frac{1}{s^2+1} = U(0)$$

The solution to the canonical form is

$$U = Ae^{-\sqrt{s}x} + Be^{\sqrt{s}x}$$

Notice that since  $U$  at  $x=\infty$  must be bounded, the  $e^{\sqrt{s}x}$  term must make the  $B$  term equal 0.

$$U = Ae^{-\sqrt{s}x} = U(0)e^{-\sqrt{s}x}$$

Apply BC at  $x=0$

$$U = \frac{1}{s^2+1} e^{-\sqrt{s}x}$$

Apply  $\mathcal{L}^{-1}$ . Find transforms in appendix.

$$1) \mathcal{L}^{-1}\left(\frac{e^{-\sqrt{s}a}}{s}\right) = \operatorname{erfc}\left(\frac{a}{2}\sqrt{t}\right)$$

$$2) \mathcal{L}^{-1}\left(\frac{s}{s^2+1}\right) = \cos(t)$$

← to get the  $\frac{1}{s}$ , mult and divide by  $s$   
and get  $\frac{s}{s^2+1}$

Apply Convolution since we have  $\mathcal{L}^{-1}$  of two parts

$$\mathcal{L}^{-1}(\mathcal{L}(f)\mathcal{L}(g)) = f * g \quad \leftarrow \text{compose, NOT "times"}$$
$$= \int_0^t f(\tau)g(t-\tau)d\tau$$

So our solution is

$$u(x,t) = \int_0^t e^{-\tau} \text{erfc}\left(\frac{x}{2\sqrt{\tau}}\right) \cos(t-\tau) d\tau$$